

H1

# PH107 Formulas

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*Joyfully crafted by Jujhaar Singh :)*

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## LM-2: Black Body Radiation

### H3 Rayleigh-Jean's Law

$$U(\nu)d\nu = \frac{8\pi\nu^2}{c^3} k_B T d\nu$$

### H3 Planck's Law

$$U(\nu)d\nu = \frac{8\pi\nu^2}{c^3} \cdot \frac{h\nu}{e^{h\nu/k_B T} - 1} d\nu = \frac{8\pi\nu^2}{c^3} k_B T \cdot \frac{h\nu/k_B T}{e^{h\nu/k_B T} - 1} d\nu$$

$$QC = \frac{h\nu/k_B T}{e^{h\nu/k_B T} - 1}, \text{ and thus we obtain:}$$

$$U(\nu)d\nu = \frac{8\pi\nu^2}{c^3} k_B T d\nu \cdot QC(\text{quantum correction})$$

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## LM-3: Compton Effect

$$\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

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## LM-4: Heat Capacity and Quantum Theory

### H3 Heat Capacity of Gases

$$c_v = \frac{f}{2} R$$

Where  $f$  is the number of degrees of freedom of the molecules of the gas

Energy available per molecule at room temperature is given by

$$E/\text{molecule} \approx 25\text{meV}$$

### H3 Dulong-Petit's Law: Classical Theory for heat capacity of solids

for a single atom moving in one direction  $\langle E \rangle = k_B T$

for a single atom moving in all 3 directions  $\langle E \rangle = 3k_B T$

$$E = 3Nk_B T = 3RT$$

$$c_v = \frac{dE}{dT} = 3R$$

### H3 Einstein's Quantum Mechanical Theory for heat capacity of Solids

$$E_n = \left( n + \frac{1}{2} \right) \cdot h\nu_E$$

for one direction:

$$E = \frac{h\nu}{e^{h\nu/k_B T} - 1}$$

and for that same atom oscillating in all 3 directions, we have:

$$E = \frac{3h\nu}{e^{h\nu/k_B T} - 1}$$

$$E = N \cdot \frac{3h\nu}{e^{h\nu/k_B T} - 1} = 3Nk_B T \cdot \frac{h\nu/k_B T}{e^{h\nu/k_B T} - 1} = 3RT \cdot \frac{h\nu/k_B T}{e^{h\nu/k_B T} - 1}$$

We define  $\theta_E = h\nu/k_B$  as the Einstein temperature of the solid and thus get

$$E = 3RT \cdot \frac{\theta_E/T}{e^{\theta_E/T} - 1}$$

In all the equations stated above,  $\nu = \nu_E$ , which is known as the Einstein frequency of the solid.

### H3 Debye Model

A few new assumptions made by Debye

$$\nu_{max} = \nu_D$$

$$\lambda_{min} = \lambda_D$$

$$\lambda_D = 2d \text{ where } d \text{ is distance between atoms}$$

At low temperatures

$$c_v \propto T^3$$

## H2 LM-5: Wave Particle duality and de Broglie's hypothesis

### H3 de Broglie Hypothesis

$$\lambda_D = \frac{h}{p}$$

### H3 Bragg's Law

$$(\text{path difference})\Delta\lambda = 2d \sin \theta$$

Where  $\theta$  is the angle of the incident rays with the surface of the lattice.

### H3 Davisson-Germer Experiment

We generally look at only first order phenomena, ie  $\Delta\lambda = \lambda$ , and find that for the angle b/w the electron gun and detector being  $\phi$  and applying Bragg's Law, we get

$$\lambda = 2d \sin \theta = 2d \cos \phi/2$$

## H2 LM-6: Wave Packets, Group Velocity and Phase Velocity

Here are some basic and useful formulas to keep in mind from here on

$$p = \hbar k = \frac{h}{\lambda}$$

$$\text{for particles: } E = \frac{p^2}{2m} = \frac{1}{2}mv^2$$

$$\text{for photons: } E = \hbar\omega = h\nu$$

### H3 Group and Phase Velocity

$$v_p = \frac{\omega}{k} = \nu\lambda$$

$$v_g = \frac{d\omega}{dk}$$

For photons, we also define

$$v_p = \frac{\omega}{k} = \frac{E}{p}$$

$$v_g = \frac{d\omega}{dk} = \frac{dE}{dp}$$

### H3 Dispersive and non-dispersive mediums

Whenever  $v_g \neq v_p$  we say that the medium is dispersive:

- if  $v_p > v_g \implies$  *normal dispersion*
- if  $v_g > v_p \implies$  *anomalous dispersion*

When we have the condition  $v_p = v_g$  we say that medium is non dispersive.

## H2 LM-7: Fourier Transform and Heisenberg's Uncertainty Principle

### H3 Heisenberg's Uncertainty Principle

1. In terms of momentum and position

$$\Delta p_k \cdot \Delta k \geq \frac{\hbar}{2} \text{ where } k \in \{x, y, z\}$$

2. In terms of energy and time

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2}$$

## H2 LM-8: The Schrödinger Equation and its properties

### H3 Schrödinger Equation

- Time Dependent Schrödinger Equation(TDSE)

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + U\Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

- Time Independent Schrödinger Equation(TISE)

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + U\Psi = E\Psi$$

Normalisation of a wave function

$$\int_{-\infty}^{\infty} |\Psi(x)|^2 dx = 1$$

### H3 Observables and Operators

Observable	Symbol	Operator
Position	$\hat{x}$	$x$
Momentum	$\hat{p}$	$-i\hbar \frac{\partial}{\partial x}$
Potential Energy	$\hat{U}$	$U(x)$
Kinetic Energy	$\hat{K}$	$\frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$
Total Energy	$\hat{E}$	$i\hbar \frac{\partial}{\partial t}$

For a normalised wave function:

$$\langle o \rangle = \int_{-\infty}^{\infty} \Psi^* \hat{O} \Psi dx$$
$$\langle o^2 \rangle = \int_{-\infty}^{\infty} \Psi^* \hat{O}^2 \Psi dx$$

### H3 Eigen functions and values

$$\hat{O}\Psi = e\Psi$$

- $\Psi$  is an Eigen function of the operator  $\hat{O}$
- $e$  is the Eigen value

## H2 Some more topics

### H3 Relativistic effects

When to consider it:

- $v$  is close to  $c$
- Energy/Kinetic Energy of electron is comparable to or larger than rest energy of electron ( $m_e c^2 = 511 \text{keV}$ )

What are the effects of considering relativity:

- $E_{total} = \sqrt{m_0^2 c^4 + p^2 c^2}$ , this includes the rest mass energy
- $KE = E_{total} - m_0 c^2$ , where  $E_{rest} = m_0 c^2$
- Now you may not use  $KE = p^2 / 2m$  and should use only the above definition

### H3 Boltzmann distribution

It states that the probability of an atom to be in a state  $i$ , ie  $p_i$ , is proportional to the given expression, where  $E_i$  is the energy of the state  $i$ :

$$p_i \propto \exp \frac{-E_i}{k_B T}$$